

LETTER

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The first stars formed soon after the plasma to gas transition at 0.3 Myr, not 200 Myr. CHG

Extremely metal-poor stars from the cosmic dawn in the bulge of the Milky Way

See journalofcosmology.com Vol 25 etc.

L. M. Howes¹, A. R. Casey², M. Asplund¹, S. C. Keller¹, D. Yong¹, D. M. Nataf¹, R. Poleski^{3,4}, K. Lind⁵, C. Kobayashi^{1,6}, C. I. Owen¹, M. Ness⁷, M. S. Bessell¹, G. S. Da Costa¹, B. P. Schmidt¹, P. Tisserand^{1,8}, A. Udalski³, M. K. Szymański³, I. Soszyński³, G. Pietrzyński^{3,9}, K. Ulaczyk^{3,10}, L. Wyrzykowski³, P. Pietrukowicz³, J. Skowron³, S. Kozłowski³ & P. Mróz³

Failure to predict star formation is one of many failures of LCDMHC cosmology vs. HGD.

The first stars are predicted to have formed within 200 million years after the Big Bang¹, initiating the cosmic dawn. A true first star has not yet been discovered, although stars^{2–4} with tiny amounts of elements heavier than helium ('metals') have been found in the outer regions ('halo') of the Milky Way. The first stars and their immediate successors should, however, preferentially be found today in the central regions ('bulges') of galaxies, because they formed in the largest over-densities that grew gravitationally with time^{5,6}. The Milky Way bulge underwent a rapid chemical enrichment during the first 1–2 billion years⁷, leading to a dearth of early, metal-poor stars^{8,9}. Here we report observations of extremely metal-poor stars in the Milky Way bulge, including one star with an iron abundance about 10,000 times lower than the solar value without noticeable carbon enhancement. We confirm that most of the metal-poor bulge stars are on tight orbits around the Galactic Centre, rather than being halo stars passing through the bulge, as expected for stars formed at redshifts greater than 15. Their chemical compositions are in general similar to typical halo stars of the same metallicity although intriguing differences exist, including lower abundances of carbon.

Stars with a low content of heavy elements have distinct spectral flux distributions, which are reflected in their colours. Using the photometric filter system on the SkyMapper telescope operated by the Australian National University, it is possible to identify metal-poor candidate stars¹⁰ in the Galactic halo⁴ and bulge⁹. We have observed ~14,000 bulge stars preselected from SkyMapper photometry using the AAOmega spectrograph on the Anglo-Australian Telescope (AAT), which enables the acquisition of 400 simultaneous stellar spectra over a 2-degree field of view. More than 500 stars with an iron abundance less than 1/100th of the solar value have been identified, making our survey the first to successfully target metal-poor stars in the Milky Way bulge. Twenty-three of these stars, targeted as the most metal-poor ones on the basis of the intermediate resolution spectra (Extended Data Table 1), were observed in June 2014 with the MIKE high-resolution spectrograph on the 6.5-m Magellan Clay telescope¹¹ to enable a comprehensive determination of their chemical compositions (Fig. 1).

The stars' effective temperatures were derived through fitting the observed hydrogen lines with theoretical spectra, while neutral and ionized iron lines provided measurements of the surface gravities and metallicities in the framework of 1D stellar atmosphere models¹² and non-equilibrium spectral line formation¹³ (Extended Data Table 2). All 23 stars were found to have $[\text{Fe}/\text{H}] \leq -2.3$, including nine stars with $[\text{Fe}/\text{H}] < -3$ (here $[\text{A}/\text{B}] = \log_{10}(N_{\text{A}}/N_{\text{B}}) - \log_{10}(N_{\text{A}}/N_{\text{B}})_{\odot}$, where $N_{\text{A}}/N_{\text{B}}$ refers to the number ratio of atoms of elements A and B in the star (* subscript) and the Sun (\odot subscript)).

The most metal-poor star, SMSS J181609.62–333218.7, has $[\text{Fe}/\text{H}] = -3.94 \pm 0.16$. The abundances of an additional 22 elements were determined spectroscopically, including the α -elements Mg, Si, Ca, and Ti, and the neutron-capture elements Y, Zr and Ba (Extended Data Table 1).

To constrain the orbits of the stars, their distances have been determined. Using the spectroscopic temperatures and surface gravities, and an assumed mass of $0.8M_{\odot}$, distances were inferred, which in nearly all cases are consistent with them being located within the bulge (Fig. 2). We have measured velocities for ten

of our stars, which confirm interesting observations that confirm hydro-gravitational-dynamics HGD cosmology vs. LCDMHC (failed) cosmology.

kinematic information are too uncertain to constrain the orbits sufficiently. Seven out of the ten stars with accurate kinematics are shown to have tightly bound orbits, placing them in the inner regions of the Milky Way (Fig. 2). In particular, using a cut-off radius of 3.43 kpc as the radius of the bulge component¹⁵, the most metal-poor star SMSS J181609.62–333218.7 has an orbit entirely contained within the bulge. Only two out of the ten stars are on much larger orbits, being merely halo stars currently passing through the bulge region. Extending these numbers to the whole sample, we can expect ~14 of the 23 bulge stars analysed here to have orbits fully within the central regions of the Milky Way; with the imminent arrival of kinematic data from the Gaia satellite, accurate orbits for all of the bulge stars will be able to be determined.

The very first stars are predicted to have brought about the cosmic dawn by forming in the centres of the largest dark matter mini-haloes, which subsequently accreted material to become the inner regions of the largest galaxies¹⁶. The typical redshift of formation for stars in the

bulge are expected to have $z \approx 5$ for halo stars. Of these, only a few are expected to have accurate orbit information. There is no cold dark matter so it cannot clump, and there are no minihalos.

$< -8 \times 10^{-4} \text{ km}^2 \text{ s}^{-2}$, which is consistent with a formation redshift of $z > 15$ (ref. 5). Low binding energies imply that the stars have been in the Galactic potential well for some time and it is very unlikely they have been accreted from a recent dwarf spheroidal merger. Their low metallicities, orbits and binding energies make these stars prime candidates for being direct descendants of the very first stars, probing a cosmic epoch otherwise completely inaccessible currently. Direct age determinations of these ancient and extremely metal-poor bulge stars from comparison with stellar evolutionary tracks or radioactive U or Th dating are currently not possible, but asteroseismic ages could possibly be inferred with the extended Kepler mission or future satellites.

¹Research School of Astronomy and Astrophysics, Australian National University, Australian Capital Territory 2601, Australia. ²Institute of Astronomy, University of Cambridge, Madingley Road, Cambridge CB3 0HA, UK. ³Warsaw University Observatory, Aleje Ujazdowskie 4, 00-478 Warszawa, Poland. ⁴Department of Astronomy, Ohio State University, 140 West 18th Avenue, Columbus, Ohio 43210, USA. ⁵Department of Physics and Astronomy, Division of Astronomy and Space Physics, Uppsala University, Box 516, SE-751 20 Uppsala, Sweden. ⁶School of Physics, Astronomy and Mathematics, Centre for Astrophysics Research, University of Hertfordshire, College Lane, Hatfield AL10 9AB, UK. ⁷Max-Planck-Institut für Astronomie, Königstuhl 17, D-69117 Heidelberg, Germany. ⁸Sorbonne Universités, UPMC Université Paris 6 et CNRS, UMR 7095, Institut d'Astrophysique de Paris, 98 bis Boulevard Arago, 75014 Paris, France. ⁹Universidad de Concepción, Departamento de Astronomía, Casilla 160-C, Concepción, Chile. ¹⁰Department of Physics, University of Warwick, Gibbet Hill Road, Coventry CV4 7AL, UK.

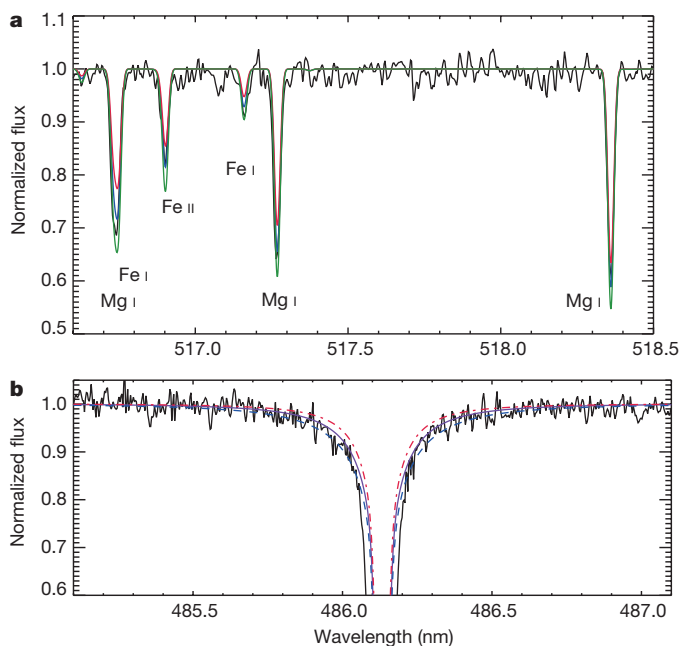


Figure 1 | Extracts of the spectrum of the lowest-metallicity star in our sample. **a**, A section of the spectrum of SMSS J181609.62–333218.7 (black line), the most metal-poor bulge star known. In blue is the predicted spectrum with the inferred stellar parameters (effective temperature $T_{\text{eff}} = 4,809$ K, $\log(g) = 1.93$ (here surface gravity g is in cgs units), $[\text{Fe}/\text{H}] = -3.94$, $[\text{Mg}/\text{Fe}] = 0.20$), and the red and green lines show spectra with all abundances scaled to ± 0.15 dex, respectively. All three predicted spectra were created using the 1D local thermodynamic equilibrium (LTE) spectrum synthesis programme, MOOG²⁶. **b**, The H β line of the same star, compared to three synthetic spectral line profiles²⁷ computed with $T_{\text{eff}} = 4,640$ K (red, dash-dot), 4,800 K (purple, continuous), and 4,960 K (blue, dashed).

Given their extremely low metallicities and large formation redshifts, these stars are likely to have formed from gas polluted by ejecta from a single or at most a few supernovae of the first stellar generation. A chemical composition analysis has been carried out to search for tell-tale nucleosynthetic signatures and possible differences from halo stars at the same metallicities. For most elements, the chemical compositions of the 23 bulge stars are consistent with typical halo stars, suggesting enrichment by similar supernovae in spite of the distinct environments and formation redshifts. Subtle differences do exist however, most notably in terms of the carbon abundances. None of the 23 stars have the large observed carbon enhancements that occur frequently in halo stars. Applying evolutionary corrections to the surface carbon abundance to counter the mixing that occurs with material processed by H-burning through CNO-cycling at late stages of the stellar lifetime¹⁷, still only one of the stars would have had a natal $[\text{C}/\text{Fe}] > 1$. In the halo, the percentage of stars that are carbon-enhanced increases dramatically at lower metallicities—from 27% of stars with $[\text{Fe}/\text{H}] < -2$ up to 69% with $[\text{Fe}/\text{H}] < -4$ (ref. 17). From the literature data on halo stars with similar iron abundances to our stars¹⁷, the probability of selecting at most one carbon-enhanced star out of 23 halo stars is only 0.2%. Carbon-enhanced stars come in two varieties: those with and those without large excesses of neutron-capture elements. The former are most likely to have been formed by mass-transfer from a binary companion that underwent the asymptotic giant branch phase. Those carbon-enhanced stars with neutron-capture excesses occur most frequently at metallicities of $[\text{Fe}/\text{H}] > -3$, whereas those without do not appear to have binary companions, and are more common at the very lowest metallicities. As none of our bulge stars are classified as having large abundances of neutron capture elements, the likelihood of finding one such carbon-enhanced star out of 23 is 7% if the frequency is the same for the bulge as for the halo. A lower

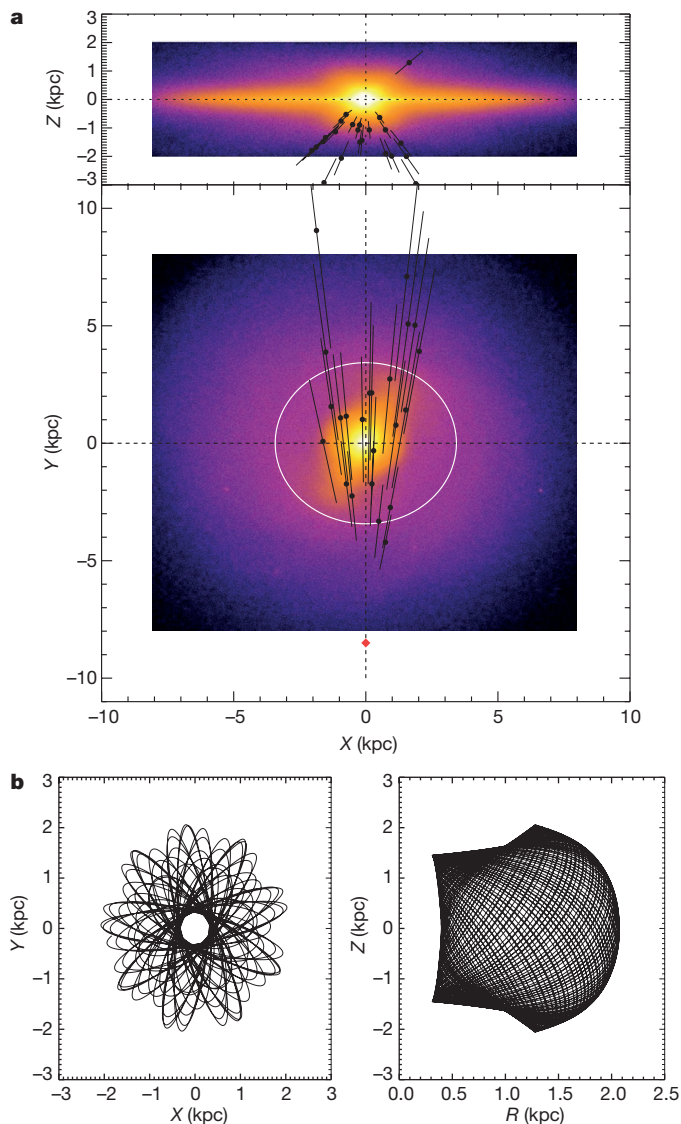


Figure 2 | The Galactic positions and orbits of the 23 stars observed at high resolution. **a**, Surface density map of a model of the Galactic bulge projected onto the X - Z (top) and X - Y (bottom) planes²⁸, where X , Y , and Z are Cartesian coordinates with the origin at the Galactic Centre and Z perpendicular to the plane of the Galaxy. Plotted over this (filled black circles) are the 23 stars of this study, with distance uncertainties shown as error bars, and a circle of radius 3.43 kpc (white: the cut-off radius of the inner bulge determined from 2MASS data¹⁵). The position of the Sun is shown with a red diamond, at 8.5 kpc from the Galactic Centre. **b**, Projections of the orbit of the lowest metallicity star, SMSS J181609.62–333218.7, both in the (R, Z) plane (right), where R is the radial direction, and in the plane of the orbit itself (left).

frequency of carbon-enhanced stars in the bulge relative to the halo is contrary to theoretical predictions; the expected dependence of the initial mass function on the cosmic microwave background¹⁸ would result in a greater number of carbon-enhanced stars near the centre of the Galaxy.

The most metal-poor bulge star, SMSS J181609.62–333218.7, is at least an order of magnitude more iron-deficient than previously found low-metallicity bulge stars^{9,19}. We have not been able to detect C in its spectrum, instead finding only an upper limit to its C abundance (Extended Data Fig. 1). This makes the upper limit on its total metallicity $[\text{Z}/\text{H}] \approx -3.8$ (total mass fraction of $\text{Z} \approx 2.1 \times 10^{-6}$), where Z represents the sum of all metals, placing it amongst the four most metal-poor stars known, along with the halo star SDSS J102915+172927 (ref. 3). The low C measured in both these stars fall below the predicted

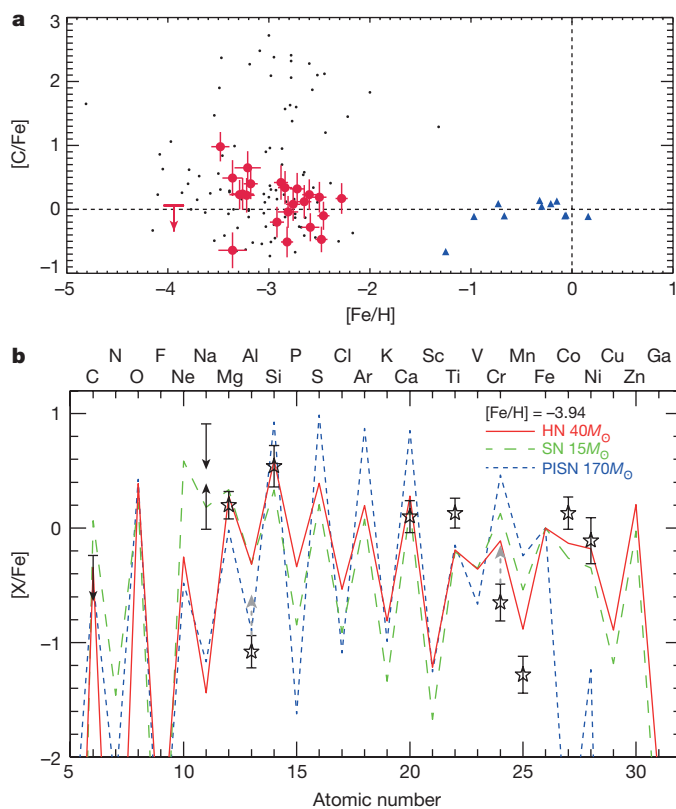


Figure 3 | Chemical abundances of the 23 stars observed at high resolution. **a**, The abundance ratio of carbon versus iron ($[C/Fe]$), with respect to metallicity ($[Fe/H]$) measured in the observed stars (filled red circles, red arrow for an upper limit). The dotted lines represent the solar abundances. Also shown for comparison are literature metal-poor halo giants (small black dots²⁴) and more metal-rich bulge stars (filled blue triangles²⁹). **b**, The chemical abundance pattern of SMSS J181609.62–333218.7, for elements X, where X is displayed at the top of the figure. Each determined abundance is shown as an open black star. These abundances are compared to three synthetic supernovae yields: a pair-instability supernova of $170M_{\odot}$ (PISN; blue, short dash²²), a core-collapse supernova of $15M_{\odot}$ (SN; green, long dash²¹), and a hypernova of $40M_{\odot}$ (HN; red, solid²¹). Dashed grey arrows represent expected non-LTE corrections; solid arrows represent measurements where only an upper or lower limit was possible. The error bars in **a** and **b** are estimates of the uncertainties in our measurements, calculated as described in Methods.

metallicity limit for formation of low-mass stars due to metal line cooling²⁰.

We have compared the detailed chemical abundance pattern of SMSS J181609.62–333218.7 to primordial supernovae yields^{21,22} (Fig. 3). In particular, the low Mg and Ca abundance, but higher Si abundance, and the absence of a pronounced odd–even abundance pattern rule out the possibility of enrichment by a pair-instability supernova resulting from a primordial star of $(140\text{--}250)M_{\odot}$. Low abundances of Cr and Mn and of α -elements, combined with the higher abundance of Co, indicate that the polluting supernova was most likely to have been a primordial hypernova—an extremely energetic kind of supernova releasing ten times the kinetic energy of regular core-collapse supernovae, possibly due to the forming black hole having larger angular momentum²³. Good agreement is found for a $40M_{\odot}$ hypernova; a more stringent Zn limit would further constrain the mass range. Unusual abundance ratios have been found in small numbers of stars in the halo—4% of halo stars with low carbon abundances have chemical peculiarities in at least two elements²⁴—but none so far appear to have been polluted by a $40M_{\odot}$ hypernova. A low $[\alpha/Fe]$ ratio (0.14 dex; here α indicates α -elements) at such low $[Fe/H]$ is consistent with an inhomogeneous enrichment from

such supernovae²⁵, while stars with higher $[\alpha/Fe]$ formed from more well-mixed gas due to a longer time delay in forming the second generation of stars.

Online Content Methods, along with any additional Extended Data display items and Source Data, are available in the online version of the paper; references unique to these sections appear only in the online paper.

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Author Contributions The project was initiated and led by M.A. The photometric target selection was made by L.M.H., C.I.O. and D.M.N. using data from the SkyMapper telescope developed by B.P.S., S.C.K., G.S.D.C., M.S.B. and P.T. The low-resolution spectra were obtained by L.M.H. and C.I.O. The data were reduced and analysed by L.M.H. using software developed by A.R.C. Target selection for the high-resolution observations was done by L.M.H., M.A. and A.R.C. with

the observations carried out by L.M.H. and D.Y.; the reduction and subsequent chemical analysis was completed by L.M.H. K.L. performed the non-LTE spectral line formation calculations, C.K. interpreted the observed chemical abundances in terms of supernova yields, and M.N. provided comparison bulge data. R.P., A.U., M.K.S., I.S., G.P., K.U., Ł.W., P.P., J.S., S.K. and P.M. obtained the OGLE observations, A.U. and M.K.S. constructed the reference images, and R.P. measured the proper motions. The manuscript was written by M.A., L.M.H. and A.R.C. with all authors contributing comments.

Author Information Reprints and permissions information is available at www.nature.com/reprints. The authors declare no competing financial interests. Readers are welcome to comment on the online version of the paper. Correspondence and requests for materials should be addressed to L.M.H. (louise.howes@anu.edu.au).

It is quite impossible to sensibly discuss the cosmic dawn starting with the nonsense cosmology LCDMHC that ignores fluid mechanics by assuming all processes of the universe are collisionless. Stars are formed as soon as the plasma turns to gas at 10^{13} seconds after the big bang. The plasma fragments to form protogalaxies at the Kolmogorov-Nomura scale of 10^{20} meters at 10^{12} seconds and the nonbaryonic dark matter begins to diffuse away because it is so weakly collisional. The missing mass of galaxies was discovered by Schild (1996) by microlensing of a quasar by a galaxy. The dark matter objects weighed $\sim 10^{24}$ kg, and formed in Jeans mass clumps of a trillion from fluid mechanical arguments of Gibson (1996). All stars are formed by binary accretion of these planets within the clumps. The clumps are called protoglobularstarclusters because when all the planets merge they are observed as globular star clusters. Infrared telescopes such as Planck and Herschel have detailed the formation of stars within PGCs. HGD cosmology is an observational fact. LCDMHC cosmology has been repeatedly falsified and belongs with speculations that the Earth is flat and that it is not necessary for doctors to wash their hands between operations.

METHODS

Observations. Photometry of the Milky Way bulge was acquired for the EMBLA survey⁹ during the commissioning period of the SkyMapper telescope in 2012 and 2013. Stars were selected from the photometry using a combination of the *g*, *i*, and *v* bandpasses, designed to give a reliable metallicity indicator¹⁰.

Spectroscopic follow-up observations took place during 2012–14, making use of the AAOmega+2dF multi-object spectrograph³⁰ on the Anglo-Australian Telescope. With between 350 and 400 stars observed in each field, spectra of more than 14,000 bulge stars have been obtained. The gratings used have a spectral resolving power of 1,300 in the blue (370–580 nm) and 10,000 in the red (840–885 nm). The data were reduced using the standard 2dfdr pipeline. Stellar spectra were fitted using a generative model that simultaneously accounts for stellar parameters (by interpolating from the AMBRE grid³¹), continuum, spectral resolution and radial velocity.

From the first two years of spectroscopic data, more than 50 stars were identified as having $[\text{Fe}/\text{H}] < -2.5$. The high-resolution spectroscopic data of 23 stars presented in this Letter are the result of observations using the MIKE spectrograph at the Magellan Clay telescope¹¹ on 15–17 June 2014. All observations were taken using a slit width of 0.7", resulting in a resolving power of 35,000 in the blue and 31,000 in the red. The data were reduced using the CarPy data reduction pipeline³², before they were normalized and summed together using the SMH software³³. The final spectra cover 330–890 nm.

Parameter and abundance determination. The stellar parameters (Extended Data Table 2) were calculated iteratively, using the original parameters from the low-resolution spectra as initial guesses. First, effective temperatures were derived by fitting the wings of the Balmer H α and H β lines with a synthetic profile (Fig. 1). These profiles were created by linearly interpolating between a grid of synthetic spectra²⁷. The best lines were fitted by a χ^2 minimization, using a weighted average of the two lines—weighting was double on the H β line, due to predicted LTE effects being larger for H α (ref. 34). The difference between the temperatures calculated for each line was on average only 26 K. The $\log(g)$, microturbulence ξ_t , and $[\text{Fe}/\text{H}]$ were then derived for that temperature, by forcing the Fe I abundance to remain constant with respect to reduced equivalent width, and equilibrium between the Fe I and Fe II abundances. Fe I and Fe II abundances were measured from the equivalent widths of a maximum of 66 Fe I lines and 24 Fe II lines (in the case of the most metal-poor star, SMSS J181609.62–333218.7, these numbers are reduced to 10 Fe I lines and 4 Fe II lines). Finally a non-LTE correction is applied to the Fe I abundance, calculated by taking the average of the line-by-line corrections¹². This correction forces an offset between the Fe I and Fe II abundances, thus replacing the initial equilibrium. This process is repeated until the parameters converge on a solution. Throughout we use the 1D MARCS model atmospheres¹³, and a shortened version of the Gaia-ESO line list, with extra lines supplemented from ref. 35 due to our wider wavelength coverage. The stellar abundances are referenced to the solar abundances of ref. 36. This analysis method was tested on seven halo stars from the literature²⁵, and the offsets found were $T_{\text{eff}} = +28$ K, $\log(g) = -0.2$, and $[\text{Fe}/\text{H}] = -0.08$ (literature values minus our values).

The abundances were measured using the equivalent widths of atomic lines (that were all on the linear part of the curve-of-growth), except in the case of C (measured from the C–H molecular bands at 431.3 nm and 432.3 nm) and Ba (synthesized in order to account for hyperfine splitting). Non-LTE corrections were calculated for Li (ref. 37), Na (ref. 38), Mg, and Ca, and applied to the individual line abundances. The literature halo abundances of Mg and Ca (ref. 25) shown in Fig. 3 have also had a NLTE correction applied, in order to ensure a fair comparison. 3σ upper limits were derived for some elements in those stars where the lines were too weak to be detected (Extended Data Table 3). The abundance offsets compared to the literature values averaged 0.10 ± 0.19 across those elements measured in common. Owing to wavelengths covered in the SkyMapper metallicity filter, it is possible that stars with extremely high C abundance appeared to be more metal-rich, and so were not selected. However, a similar study of metal-poor stars discovered in the halo with SkyMapper³⁹ found the fraction of C-enhanced stars was identical to that reported in previous surveys⁴⁰. Furthermore, we followed up 14,000 stars with intermediate resolution spectra, and determined metallicities using those spectra. The majority of the stars observed had $[\text{Fe}/\text{H}] \approx -1.0$ and included some that had solar metallicities, so it is highly unlikely that we missed any C-enhanced extremely metal-poor stars in our selection.

The systematic uncertainties in the temperature determinations were estimated to be ± 100 K, and the statistical uncertainties averaged ± 125 K, so when these are combined in quadrature we conclude the total uncertainty to be ± 160 K. The ξ_t uncertainties are estimated to be ± 0.2 , mostly due to systematics.

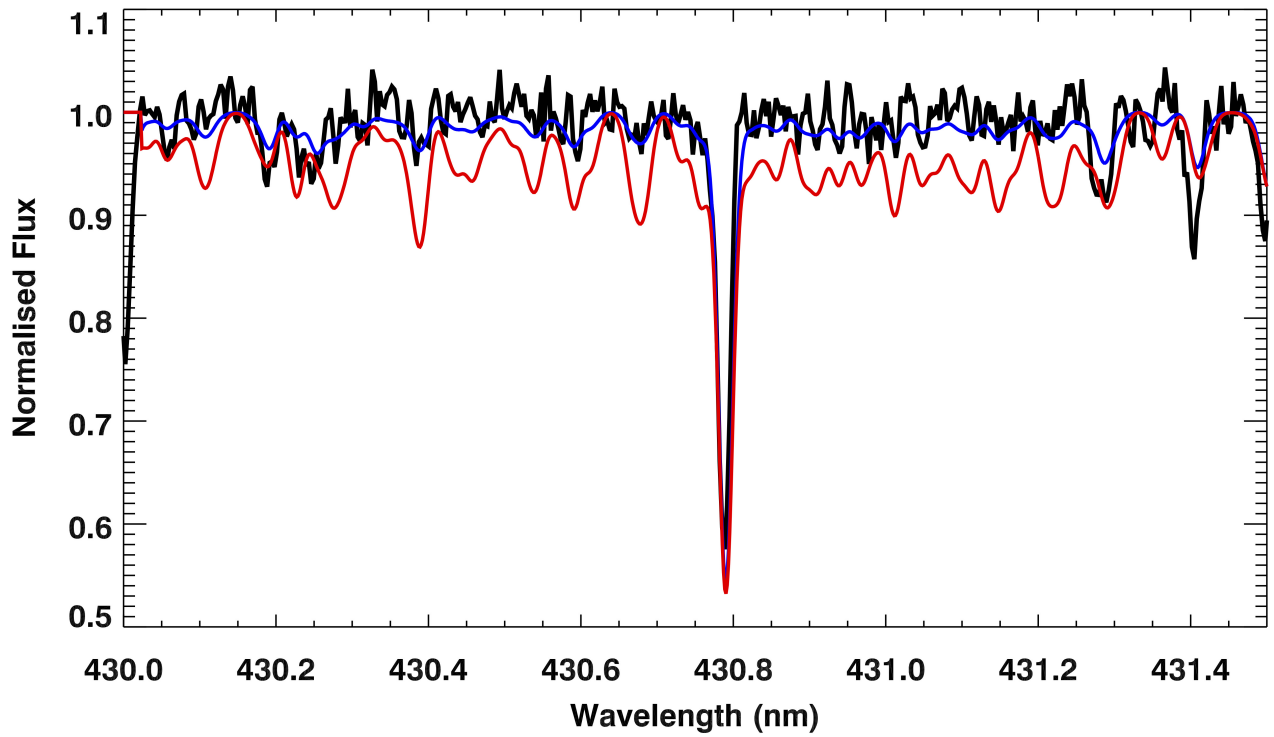
The standard errors of the individual line abundances of Fe I and Fe II were combined in quadrature to evaluate the $\log(g)$ uncertainties. The differences between the $[\text{Fe}/\text{H}]$ values when varying the temperature, surface gravity, and microturbulence by their respective errors were combined in quadrature with the standard error of the Fe II lines to produce the $[\text{Fe}/\text{H}]$ uncertainties. The individual abundance errors were also calculated using this method, using the standard error of the individual abundances for the lines of that particular element.

For SMSS J181609.62–333218.7, which has $[\text{Fe}/\text{H}] = -3.94$, a measurement of Na was not possible, owing to the 818.3 nm and 819.4 nm lines being too weak, and the Na D lines (588.9 nm and 589.5 nm) being partly blended with interstellar Na lines. We have derived a range of possible values for this star, taking the upper limit from the non-detection at 819.4 nm, and the lower limit from fitting a Gaussian to the Na D lines, taking into account the interstellar Na.

Distances and orbital parameters. Distances to the stars were calculated by comparing the absolute and apparent bolometric magnitudes. The absolute magnitudes were recovered from the relation $M_* = M_\odot - 2.5 \log(L_*/L_\odot)$, where the luminosities are calculated using $L_*/L_\odot = 4\pi\sigma T^4 M_* G / 10^{\log(g_*)}$, taking $M_* = (0.8 \pm 0.2)M_\odot$ for all stars. The apparent bolometric magnitudes are reconstructed from the 2MASS *JHK_s* magnitudes (Extended Data Table 1), assuming reddening⁴¹ (as no more-recent reddening catalogue covers all 23 stars), via the methodology of ref. 42. The proper motions are based on I band images taken during the OGLE-IV¹⁴ observations of the Galactic bulge. Relative proper motions were derived from multiple epochs of data for each field⁴³, and the uncertainties are a combination of statistical and systematic (for each star, the systematic uncertainty is estimated to be ~ 0.4 mas yr⁻¹). These were converted into absolute proper motions by adding the predicted average bulge motion for each field, calculated using the Besancon Galaxy model⁴⁴. The orbits were calculated using the python package galpy⁴⁵ and the galactic potential assumed in these calculations was a three-component Milky Way-like potential⁴⁶. To model the uncertainty distributions, we sampled 1,000 orbits using a Monte Carlo simulation, assuming a normal distribution for the uncertainties of the input parameters. The results of this are included in Extended Data Table 6. One star, SMSS J175455.52–380339.3, has an unbound E_{tot} and impractically large orbital parameters, suggesting that one or more of our input parameters need to be changed.

Code availability. All codes used to analyse the data presented are publicly available. In particular, the 1D LTE analysis used was made possible with the line analysis and spectrum synthesis code MOOG²⁶.

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Extended Data Figure 1 | The C–H band of SMSS J181609.62 – 333218.7. The C–H band is used to derive an upper limit for C in our most metal-poor star, SMSS J181609.62 – 333218.7. Synthetic spectra with abundances of $[C/Fe] = 0.06$ (blue) and $[C/Fe] = 0.56$ (red) are shown for comparison.

Extended Data Table 1 | Coordinates and 2MASS photometry of the 23 stars observed

Name (SMSS)	RA (°)	Dec (°)	l (°)	b (°)	J (mag)	H (mag)	K_S (mag)
J173823.38-145701.1	264.597	-14.950	11.1	8.7	10.85	10.22	10.03
J182048.26-273329.2	275.201	-27.558	5.0	-6.1	12.94	12.42	12.25
J183744.90-280831.1	279.437	-28.142	6.2	-9.7	12.29	11.69	11.53
J183647.89-274333.1	279.200	-27.726	6.5	-9.3	10.68	10.03	9.77
J183812.72-270746.3	279.553	-27.130	7.1	-9.3	13.38	12.79	12.61
J183719.09-262725.0	279.330	-26.457	7.7	-8.9	12.79	12.19	12.03
J184201.19-302159.6	280.505	-30.367	4.5	-11.5	14.52	14.08	14.00
J184656.07-292351.5	281.734	-29.398	5.9	-12.0	13.12	12.61	12.51
J181406.68-313106.1	273.528	-31.518	0.8	-6.6	12.12	11.56	11.35
J181317.69-343801.9	273.324	-34.634	357.9	-7.9	13.09	12.55	12.50
J181219.68-343726.4	273.082	-34.624	357.9	-7.7	12.80	12.28	12.15
J181609.62-333218.7	274.040	-33.539	359.2	-7.9	13.39	12.84	12.71
J181634.60-340342.5	274.144	-34.062	358.8	-8.3	12.56	11.99	11.90
J175544.54-392700.9	268.936	-39.450	352.0	-7.1	13.71	13.19	13.09
J175455.52-380339.3	268.731	-38.061	353.1	-6.3	11.98	11.39	11.26
J175746.58-384750.0	269.444	-38.797	352.8	-7.2	13.09	12.60	12.51
J181736.59-391303.3	274.402	-39.218	354.2	-10.8	12.06	11.54	11.37
J181505.16-385514.9	273.772	-38.921	354.2	-10.2	13.63	13.15	13.09
J181921.64-381429.0	274.840	-38.241	355.2	-10.6	13.64	13.13	13.03
J175722.68-411731.8	269.345	-41.292	350.5	-8.3	13.85	13.25	13.21
J175021.86-414627.1	267.591	-41.774	349.4	-7.4	11.74	11.23	11.17
J175636.59-403545.9	269.152	-40.596	351.1	-7.9	12.86	12.29	12.19
J175433.19-411048.9	268.638	-41.180	350.4	-7.8	11.94	11.43	11.33

RA, right ascension; Dec., declination; l and b , Galactic longitude and latitude, respectively.

Extended Data Table 2 | Stellar parameters of the 23 stars observed

Name (SMSS)	V_{helio} (km s ⁻¹)	d_{\odot} (kpc)	T_{eff} (K)	$\log g$ (cgs)	[Fe/H] (dex)	ξ_t (km s ⁻¹)	[α /Fe] (dex)
J173823.38-145701.1	46.1	8.5	4599	0.99	-3.36	2.30	0.12
J182048.26-273329.2	51.5	6.0	4949	2.22	-3.48	1.90	0.37
J183744.90-280831.1	-132.6	17.6	4597	0.98	-2.92	2.05	0.33
J183647.89-274333.1	-381.4	6.6	4649	1.17	-2.48	2.50	0.30
J183812.72-270746.3	155.3	12.3	4873	1.74	-3.22	1.81	-0.01
J183719.09-262725.0	-244.7	10.0	4791	1.64	-3.18	1.81	0.32
J184201.19-302159.6	171.8	9.6	5136	2.55	-2.84	1.96	0.30
J184656.07-292351.5	91.0	9.5	4857	1.93	-2.76	1.83	0.34
J181406.68-313106.1	4.9	9.3	4821	1.48	-2.82	1.96	0.22
J181317.69-343801.9	139.3	6.5	5015	2.25	-2.28	1.48	0.41
J181219.68-343726.4	-386.2	8.0	4873	1.94	-2.50	1.93	0.32
J181609.62-333218.7	27.4	10.4	4809	1.93	-3.94	1.60	0.14
J181634.60-340342.5	-170.3	10.5	4821	1.61	-2.46	1.79	0.06
J175544.54-392700.9	-279.6	13.5	4857	1.83	-2.65	1.60	0.32
J175455.52-380339.3	23.5	13.5	4714	1.10	-3.36	1.80	0.08
J175746.58-384750.0	-59.4	9.1	5064	1.96	-2.81	2.36	0.29
J181736.59-391303.3	-177.9	15.7	4612	1.05	-2.59	2.09	0.32
J181505.16-385514.9	202.1	5.0	4962	2.73	-3.29	2.10	0.35
J181921.64-381429.0	-97.7	11.2	4917	2.02	-2.72	1.94	0.30
J175722.68-411731.8	63.8	12.4	4894	1.97	-2.88	2.02	0.19
J175021.86-414627.1	181.4	4.1	5015	2.12	-2.60	1.55	0.30
J175636.59-403545.9	-28.8	9.8	4934	1.79	-3.21	1.96	0.20
J175433.19-411048.9	-229.3	5.6	4912	1.91	-3.26	1.94	0.35

Symbols: V_{helio} , heliocentric velocity; d_{\odot} , distance from the Sun to the star; T_{eff} , effective temperature; $\log(g)$, stellar surface gravity; ξ_t , microturbulence; [α /Fe] = ([Mg/Fe] + [Ca/Fe] + [Ti/Fe])/3. Average uncertainties: velocity, 1.0 km s⁻¹; distance, 3.0 kpc; temperature, 160 K; microturbulence, 0.2 dex; $\log(g)$, 0.14 dex; [Fe/H], 0.09 dex; [α /Fe], 0.13 dex.

Extended Data Table 3 | Chemical abundances measured for each star, C to Ca

Name (SMSS)	A(Li)	[C/Fe]	[Na/Fe]	[Mg/Fe]	[Al/Fe]	[Si/Fe]	[K/Fe]	[Ca/Fe]
J173823.38-145701.1		0.49	0.04	0.17	-0.78	0.27		0.12
J182048.26-273329.2		0.98	-0.28	0.54	-0.63	0.96		0.30
J183744.90-280831.1	0.16	-0.20	-0.28	0.44	-0.52	0.58	0.36	0.25
J183647.89-274333.1		-0.47	-0.24	0.33	-0.66	0.51		0.18
J183812.72-270746.3	0.93	0.22	-0.39	0.05	-1.23	0.14		0.03
J183719.09-262725.0		0.40	-0.19	0.47	-0.77	0.36	0.41	0.25
J184201.19-302159.6		0.34	-0.38	0.26	-0.89	0.38	0.53	0.37
J184656.07-292351.5	1.04	0.08	-0.30	0.41	-0.95	0.36	0.58	0.28
J181406.68-313106.1		-0.51	0.18	0.23	-0.94	0.32		0.16
J181317.69-343801.9	1.05	0.17	-0.33	0.53	-0.82	0.33	0.63	0.34
J181219.68-343726.4	1.01	0.19	-0.22	0.30	-0.86	0.25		0.31
J181609.62-333218.7		<0.06	-0.01<0.91 ^a	0.20	-1.08	0.54		0.00
J181634.60-340342.5		-0.10	-0.53	0.05	-1.08	0.11	0.21	0.03
J175544.54-392700.9	0.87	0.12	-0.32	0.29	-0.88	0.34	0.36	0.29
J175455.52-380339.3		-0.64		0.06	-0.88	0.30		0.03
J175746.58-384750.0		-0.04		0.37	-1.10	0.44		0.24
J181736.59-391303.3		-0.28	-0.11	0.38	-0.69	0.53	0.53	0.26
J181505.16-385514.9		0.23	-0.23	0.21	-0.96	0.15		0.36
J181921.64-381429.0	1.04	0.32	-0.24	0.28	-0.82	0.54	0.44	0.26
J175722.68-411731.8		0.42	-0.42	0.21	-0.70	0.48	0.14	0.13
J175021.86-414627.1	0.98	0.23	-0.37	0.30	-0.82	0.42		0.28
J175636.59-403545.9	0.93	0.65		0.30	-0.76	0.45		0.11
J175433.19-411048.9	0.92	0.24	-0.03	0.40	-0.74	0.43	0.44	0.32

A(Li) is the logarithmic abundance of lithium. All abundances are derived using LTE, except for Li, Na, Mg, and Ca, where non-LTE corrections have been applied. Average uncertainties: Li, 0.20; C, 0.25; Na, 0.20; Mg, 0.16; Al, 0.22; Si, 0.21; K, 0.17; Ca, 0.12.

^a-0.01 is the lower limit, and 0.91 is the upper limit; see Methods for details.

Extended Data Table 4 | Chemical abundances measured for each star, Sc to Cu

Name (SMSS)	[Sc/Fe]	[Ti/Fe]	[Cr/Fe]	[Mn/Fe]	[Co/Fe]	[Ni/Fe]	[Cu/Fe]
J173823.38-145701.1		-0.09	-0.22	-0.80	0.17	-0.21	<0.96
J182048.26-273329.2		0.16	-0.51	-0.97	0.24	-0.33	<1.33
J183744.90-280831.1	0.04	0.20	-0.23	-0.38	0.36	0.14	<0.29
J183647.89-274333.1	0.14	0.34	-0.27	-0.35	0.01	0.02	-0.43
J183812.72-270746.3		-0.20	-0.51	-0.32	0.22	-0.08	<1.06
J183719.09-262725.0	0.18	0.14	-0.33	-0.34	0.23	0.23	<1.10
J184201.19-302159.6	-0.03	0.20	-0.24	-0.57	0.35	-0.02	<0.72
J184656.07-292351.5	0.11	0.28	-0.19	-0.31	0.11	0.07	<0.45
J181406.68-313106.1	0.08	0.19	-0.30	-0.60	0.22	0.09	<0.50
J181317.69-343801.9	0.11	0.34	-0.19	-0.08	0.12	0.10	<0.13
J181219.68-343726.4	0.18	0.31	-0.15	-0.14	0.22	0.22	<0.17
J181609.62-333218.7		0.13	-0.65	-1.28	0.13	-0.11	<1.57
J181634.60-340342.5	-0.26	0.04	-0.24	-0.35	-0.20	-0.03	<-0.05
J175544.54-392700.9	-0.05	0.32	-0.22	-0.18	0.24	-0.05	<0.32
J175455.52-380339.3		0.02	-0.47	-1.05	0.07	-0.12	<0.85
J175746.58-384750.0	0.12	0.19	-0.44	-0.59	0.09	-0.01	<0.79
J181736.59-391303.3	0.15	0.24	-0.31	-0.34	0.00	0.08	<0.10
J181505.16-385514.9	0.40	0.38	-0.54	-0.81	0.17	-0.11	<1.02
J181921.64-381429.0	0.01	0.32	-0.17	-0.44	0.33	0.25	<0.50
J175722.68-411731.8	-0.22	0.15	-0.31	-0.48	-0.04	0.04	<0.82
J175021.86-414627.1	-0.23	0.26	-0.27	-0.31	0.16	0.16	<0.27
J175636.59-403545.9	-0.28	0.06	-0.29	-0.72	0.17	-0.26	<0.95
J175433.19-411048.9		0.22	-0.52	-0.91	0.29	0.02	<1.08

All abundances are derived using LTE. Average uncertainties: Sc, 0.10; Ti, 0.10; Cr, 0.21; Mn, 0.25; Co, 0.23; Ni, 0.19; Cu, 0.25.

Extended Data Table 5 | Chemical abundances measured for each star, Zn to Eu

Name (SMSS)	[Zn/Fe]	[Sr/Fe]	[Y/Fe]	[Zr/Fe]	[Ba/Fe]	[La/Fe]	[Eu/Fe]
J173823.38-145701.1	0.66	0.03	0.02	0.23	-0.04	-0.10	
J182048.26-273329.2	<1.01	-0.47			0.03	<1.33	
J183744.90-280831.1	0.27	-0.29	-0.32	0.03	-0.31	<0.12	
J183647.89-274333.1	0.23	0.18	-0.20	0.45	0.13	0.17	0.82
J183812.72-270746.3	<0.79	-1.03			-0.70	<0.77	
J183719.09-262725.0	0.48	0.04	0.53		-0.51	<1.03	
J184201.19-302159.6	<1.15	-0.20	0.04	0.70	0.16	<0.94	
J184656.07-292351.5	0.48	-0.26	-0.51	0.14	-0.32	<0.36	
J181406.68-313106.1	0.42	-1.61			-0.72	<0.32	
J181317.69-343801.9	0.17	0.17	-0.11	0.34	0.22	-0.09	0.15
J181219.68-343726.4	0.33	-0.06	-0.11	0.09	0.13	<0.90	0.48
J181609.62-333218.7	<1.40	-0.85	0.23		<-0.66	<1.09	0.91
J181634.60-340342.5	0.21	-0.25	-0.65	-0.21	-0.32	-0.14	-0.11
J175544.54-392700.9	0.36	-0.10	-0.2	0.38	-0.11	-0.15	0.21
J175455.52-380339.3	0.63	0.47	0.01	0.14	-0.57	<0.66	
J175746.58-384750.0		-0.21	0.04	0.91	0.23	<1.26	0.65
J181736.59-391303.3	0.23	-0.14	-0.47	0.14	-0.28	<1.19	0.21
J181505.16-385514.9	<0.95	-0.19	0.14	0.71	0.04	<0.54	0.96
J181921.64-381429.0	0.42	-0.21	-0.14	0.51	-0.01	0.48	0.59
J175722.68-411731.8	<0.95	-0.30	-0.30	0.24	-0.19	<0.63	0.52
J175021.86-414627.1	0.41	-0.14	-0.40	0.25	-0.08	<0.60	0.23
J175636.59-403545.9	0.86	0.55	0.24	0.71	-0.95	<1.21	
J175433.19-411048.9	0.50	-0.81	-0.29		-0.42	<0.91	

All abundances are derived using LTE. Average uncertainties: Zn, 0.10; Sr, 0.20; Y, 0.12; Zr, 0.12; Ba, 0.17; La, 0.15; Eu, 0.16.

Extended Data Table 6 | Orbital parameters

Name (SMSS)	$\mu_\alpha \cos \delta$ (mas yr ⁻¹)	μ_δ (mas yr ⁻¹)	Mean r_{peri} (kpc)	Mean r_{ap} (kpc)	Mean Eccentricity	Mean Z_{max} (kpc)	E_{tot} (10 ⁴ km ² s ²)
J182048.26-273329.2	-4.10 ±0.52	-6.38 ±0.51	0.5 ^{+0.9} _{-0.2}	2.9 ^{+1.7} _{-1.1}	0.70 ^{+0.11} _{-0.23}	0.9 ^{+0.2} _{-0.5}	-11.3 ^{+2.2} _{-1.4}
J184201.19-302159.6	-0.38 ±0.90	-0.82 ±0.90	1.2 ^{+1.1} _{-0.7}	6.6 ^{+4.3} _{-1.6}	0.72 ^{+0.13} _{-0.20}	4.3 ^{+1.7} _{-1.8}	-7.1 ^{+2.1} _{-1.2}
J184656.07-292351.5	1.17 ±0.89	-2.32 ±0.89	1.2 ^{+1.5} _{-0.8}	4.8 ^{+2.7} _{-1.4}	0.65 ^{+0.19} _{-0.41}	3.2 ^{+1.2} _{-1.4}	-8.3 ^{+1.9} _{-1.2}
J181406.68-313106.1	2.28 ±0.52	-8.25 ±0.52	1.1 ^{+2.3} _{-0.8}	3.5 ^{+9.1} _{-1.3}	0.63 ^{+0.18} _{-0.16}	2.7 ^{+8.6} _{-0.8}	-9.5 ^{+5.4} _{-1.8}
J181219.68-343726.4	-2.42 ±1.14	-1.29 ±1.14	0.7 ^{+0.8} _{-0.4}	17.3 ^{+11.3} _{-4.9}	0.92 ^{+0.05} _{-0.03}	7.2 ^{+3.2} _{-1.8}	-3.4 ^{+1.9} _{-1.4}
J181609.62-333218.7	-4.14 ±0.64	-3.74 ±0.64	1.0 ^{+3.5} _{-0.8}	3.4 ^{+3.3} _{-1.8}	0.53 ^{+0.23} _{-0.28}	1.9 ^{+2.6} _{-0.9}	-9.9 ^{+4.0} _{-2.6}
J181634.60-340342.5	1.92 ±0.62	-0.31 ±0.62	1.9 ^{+2.0} _{-1.3}	7.9 ^{+8.3} _{-2.4}	0.65 ^{+0.15} _{-0.10}	3.9 ^{+4.2} _{-1.7}	-6.3 ^{+3.0} _{-2.0}
J175544.54-392700.9	0.03 ±1.49	-0.35 ±1.46	1.7 ^{+2.3} _{-1.2}	29.3 ^{+67.4} _{-16.5}	0.92 ^{+0.05} _{-0.04}	10.8 ^{+29.6} _{-6.0}	-1.4 ^{+3.7} _{-3.2}
J175455.52-380339.3	1.98 ±1.14	4.76 ±1.14	5.5 ^{+4.8} _{-3.9}	811.5 ^{+1216.1} _{-796.2}	0.98 ^{+0.01} _{-0.16}	81.0 ^{+259.4} _{-76.6}	6.9 ^{+12.2} _{-10.7}
J175746.58-384750.0	1.86 ±1.25	0.17 ±1.25	1.8 ^{+1.9} _{-1.2}	5.9 ^{+5.0} _{-1.9}	0.59 ^{+0.24} _{-0.28}	2.3 ^{+2.8} _{-1.1}	-7.5 ^{+2.9} _{-1.8}

Symbols: $\mu_\alpha \cos \delta$ and μ_δ are the proper motions in equatorial coordinates; r_{peri} and r_{ap} are the pericentric and apocentric radii of the orbit, respectively, Z_{max} is the maximum distance the orbit reaches above/below the Galactic plane, and E_{tot} is the total energy of the orbit. All values given here are the mean values from the Monte Carlo simulation of 1,000 orbits.