

On intense vortex structures in isotropic turbulence

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Johns Hopkins web-based data base
(<http://turbulence.pha.jhu.edu>)
DNS forced homogeneous isotropic
turbulence
 1024^3 $Re_\lambda = 433$

Intense vortex structures

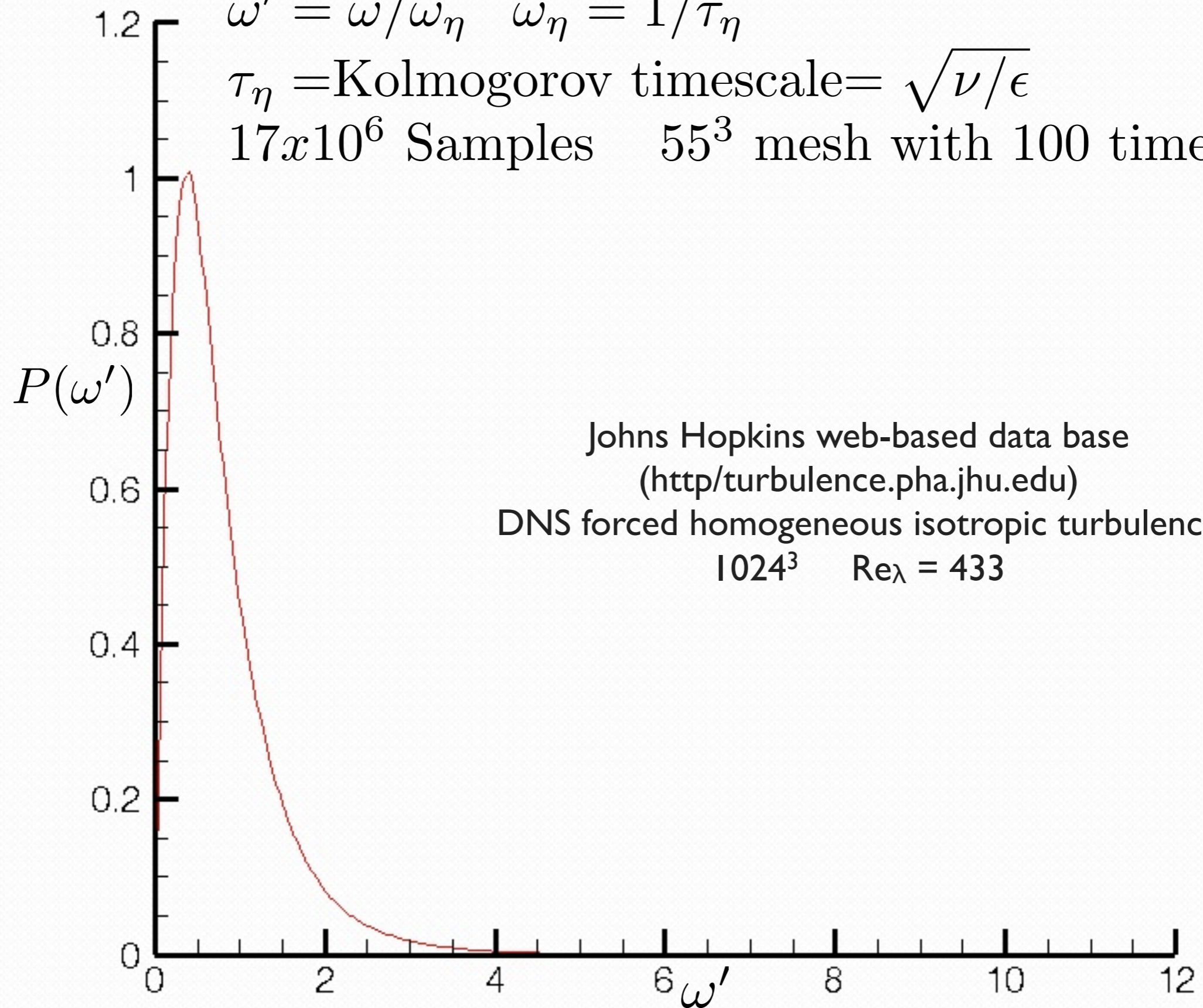
- PDF of vorticity amplitudes - asymptotics
- Geometry of intense structures, vorticity distribution
- Evolution in time

PDF of Normalized Vorticity Amplitudes

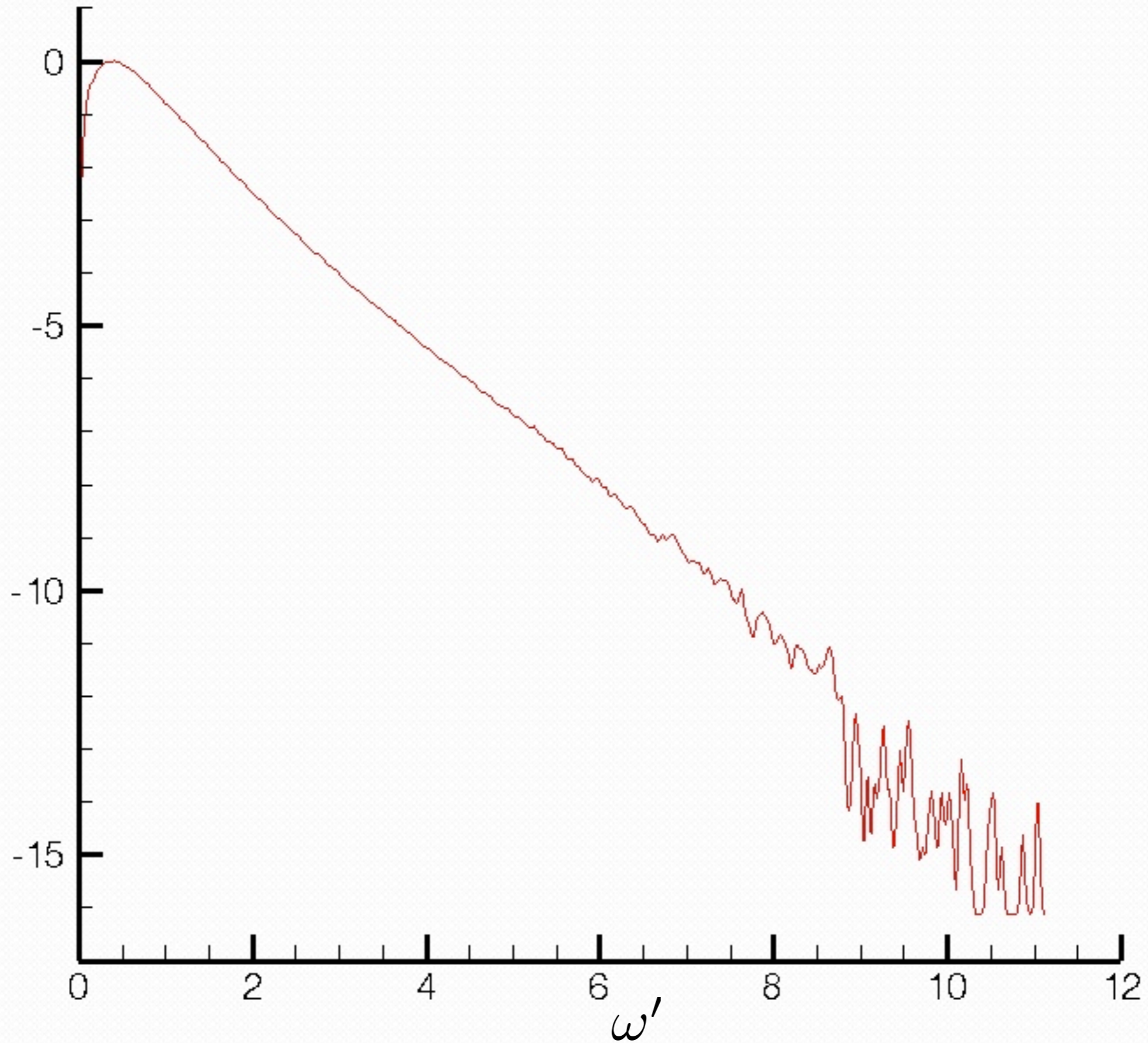
$$\omega' = \omega / \omega_\eta \quad \omega_\eta = 1 / \tau_\eta$$

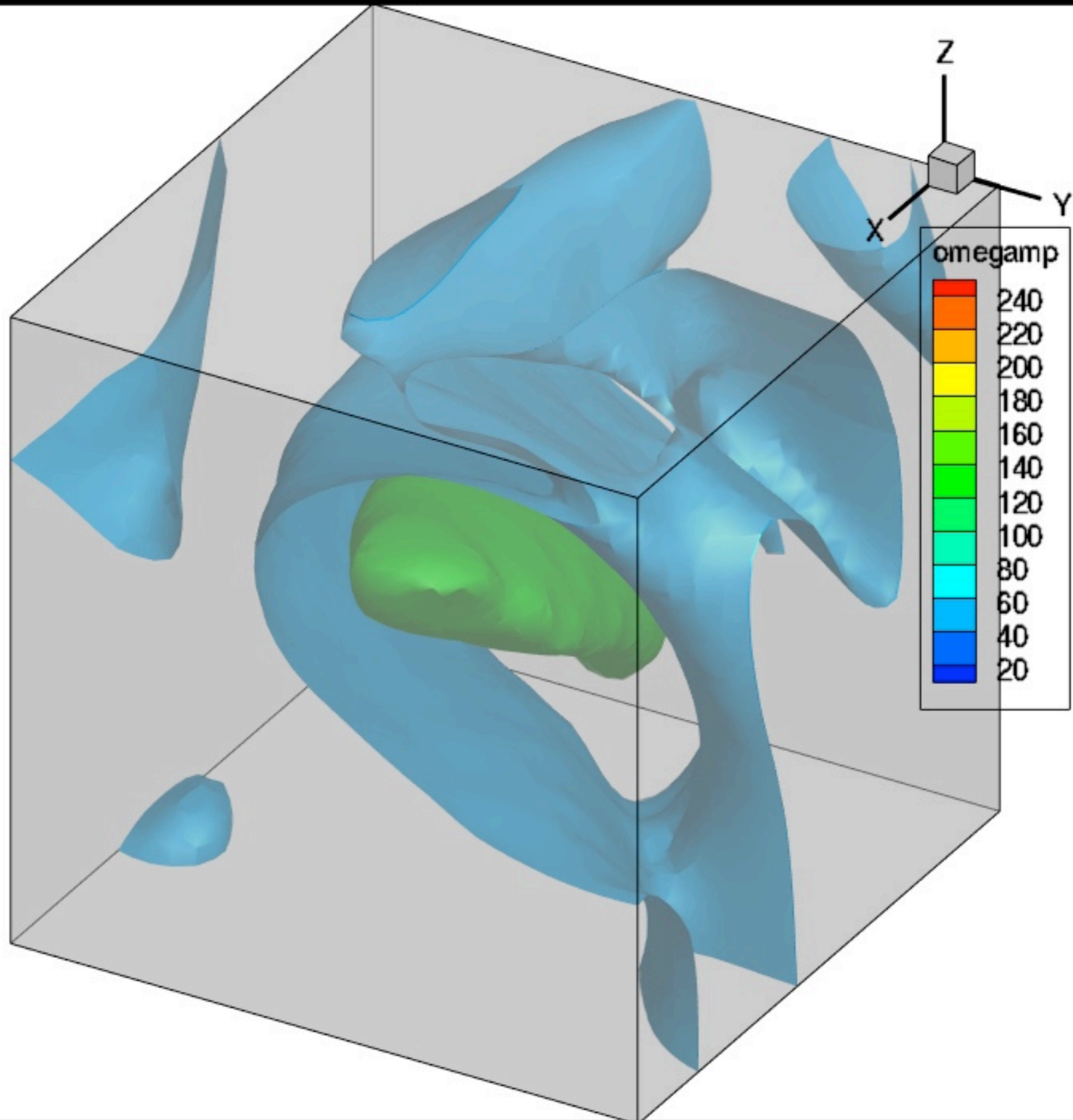
$$\tau_\eta = \text{Kolmogorov timescale} = \sqrt{\nu / \epsilon}$$

17×10^6 Samples 55^3 mesh with 100 timesteps

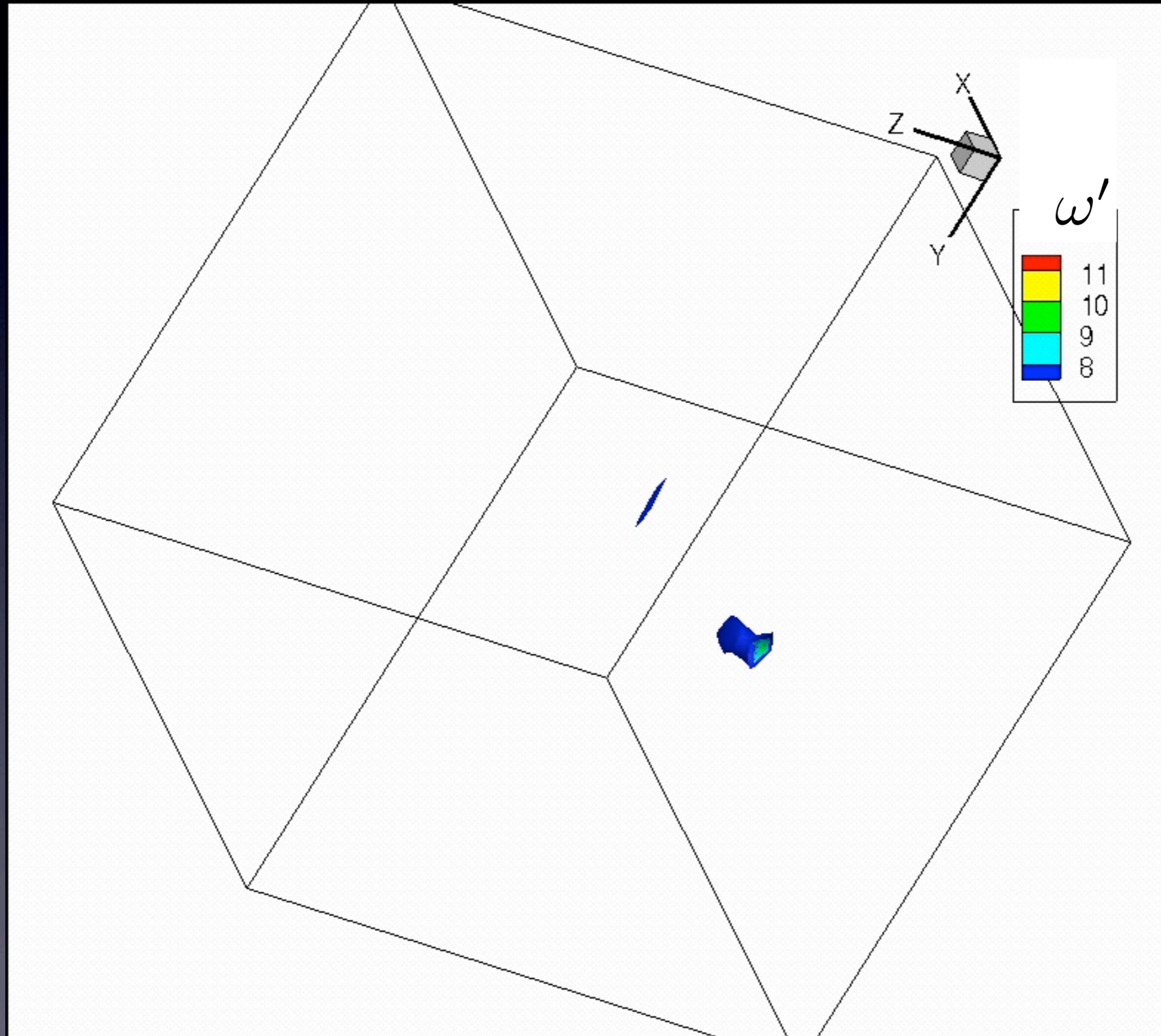


$\text{Log}P(\omega')$

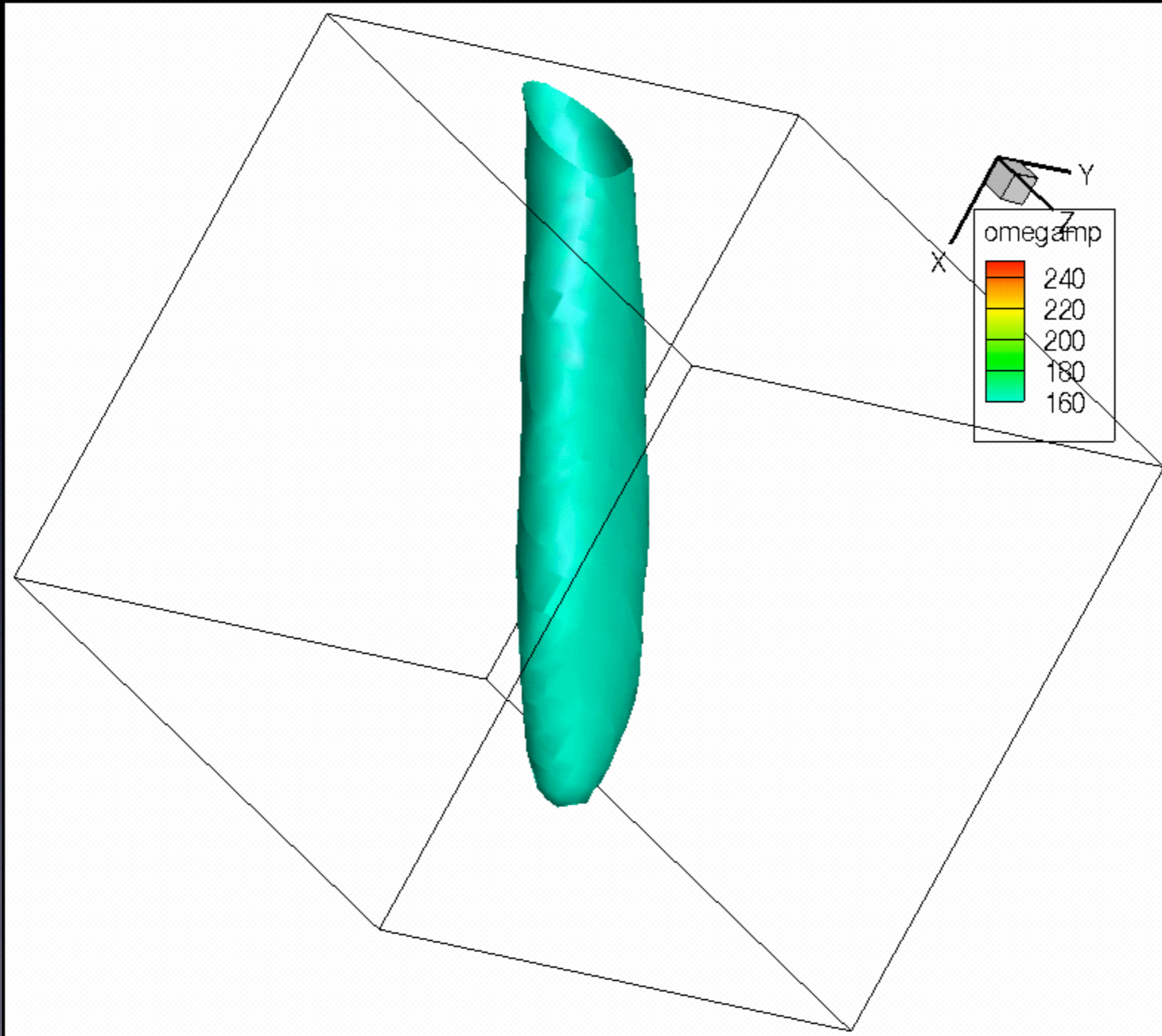


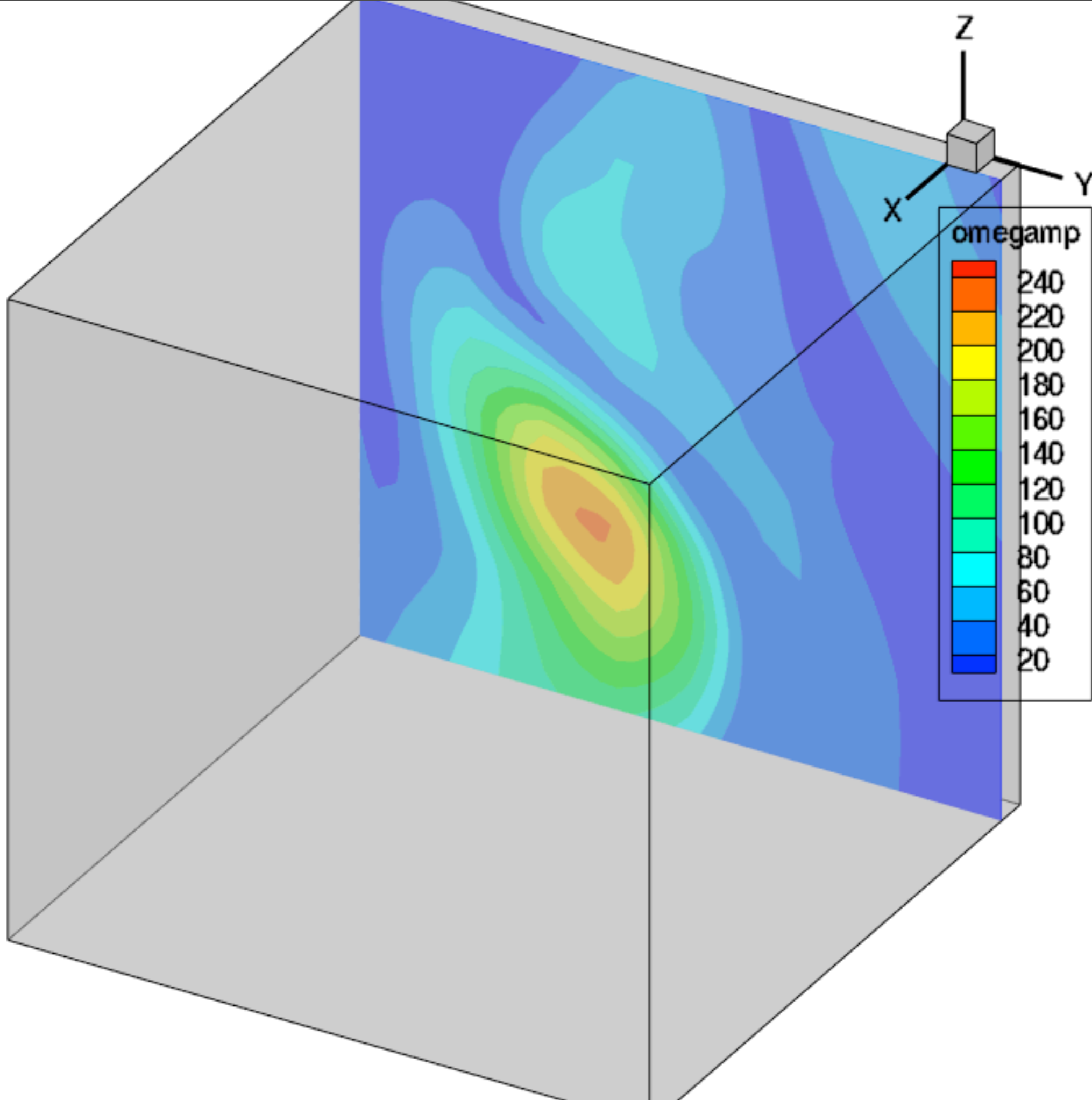


High Amplitude ω'
Box size $(32\Delta)^3 = (68\eta)^3 = (1.7\lambda)^3$



Intense ω' Region (zoom in)
Box size $(9\Delta)^3 = (19\eta)^3 = (0.47\lambda)^3$





PDF for a single structure - $P_S(\omega)$

Model I - Exponential decay of circulation

$$\omega_x = \frac{\Gamma}{\pi\sigma_r^2} \exp(-r^2/\sigma_r^2) \exp(-x^2/\sigma_x^2)$$

$$\text{Let } S = r^2/\sigma_r^2 + x^2/\sigma_x^2 \rightarrow S = \log\left(\frac{\omega_{max}}{\omega}\right)$$

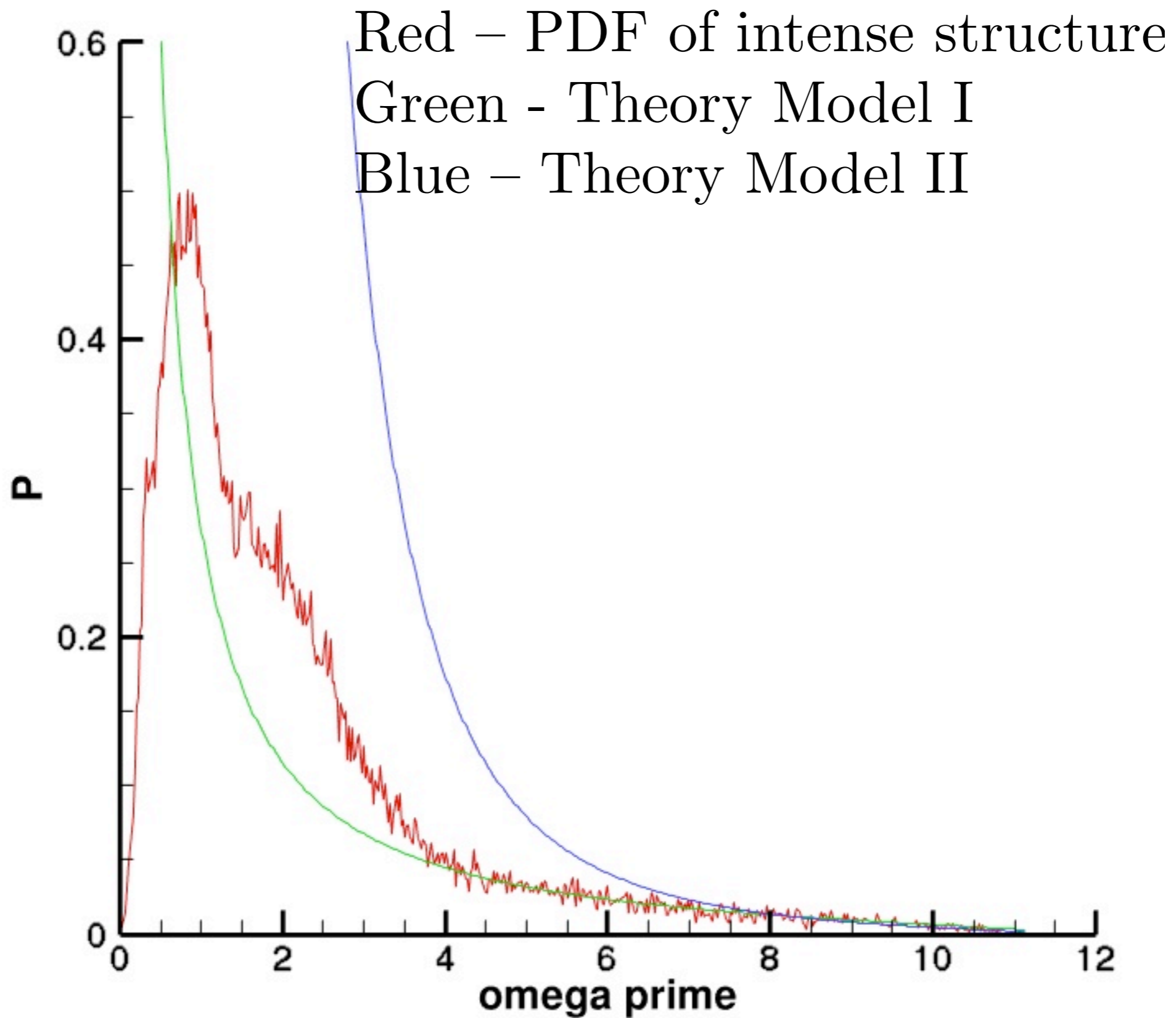
$$\text{Volume inclosed by } S \rightarrow V(S) = \frac{4}{3}\pi\sigma_r^2\sigma_x S^{3/2}$$

$$P_V(V) = \frac{dV}{V_T} \rightarrow \boxed{P_S(\omega)d\omega = C \left[\log\left(\frac{\omega_{max}}{\omega}\right) \right]^{1/2} \frac{d\omega}{\omega}}$$

Model II - Constant circulation, expanding vortex core

$$\omega_x = \frac{\Gamma}{\pi\sigma_r^2(x)} \exp(-r^2/\sigma_r^2(x))$$

$$\text{with } \sigma_r^2(x) = \sigma_o^2(1 + x^2/\sigma_x^2)$$



Structure evolution in time (preliminary)

Assume the following form for ω_x

$$\omega_x = \frac{\Gamma(t)}{\pi\sigma_r^2(t)} \exp(-r^2/\sigma_r^2(t)) \exp(-x^2/\sigma_x^2(t))$$

Consider the vorticity transport equation for ω_x

with leading order straining flow -

$$\mathbf{U} = (a(t)x + b(t)(x^2 - r^2/2) + c(t)(x^3/3 - xr^2/2))\mathbf{e}_x \\ + (-a(t)r/2 - b(t)xr - c(t)(x^2/2 - r^3/8))\mathbf{e}_r$$

Exact equation for $\Gamma(t)$ (Kelvin's Theorem)

$$\frac{d\Gamma(t)}{dt} = -2\nu/\sigma_x^2(t)$$

Structure evolution in time (continued, preliminary)

Compute $\int dV$, and $\int r^2 dV$ moments of the transport equation with the result

$$\frac{d\sigma_x^2(t)}{dt} = 2a(t)\sigma_x^2(t) + c(t)\sigma_x^4(t) + 4\nu$$

$$\frac{d\sigma_r^2(t)}{dt} = -a(t)\sigma_r^2(t) - \frac{c(t)}{2}\sigma_r^4(t) + 4\nu$$

with ω_r term and local induced velocities argued to be small. The straining flow proportional to $b(t)$ simply shifts the location of the maximum ω_x along the x axis.

Quasi-steady solution for the geometry if $c < 0$

$$\sigma_x^2(t) \approx -2\frac{a(t)}{c(t)}$$

$$\sigma_r^2(t) \approx 2\frac{\nu}{a}$$

$\Gamma(t)$ decays exponentially as given above.

Summary

Intense vortex structures

- Geometry of intense structures, vorticity distribution
Cigar shaped, aspect ratio ≈ 6.5 , gaussian distribution of circulation along axis
- Connection between geometry and PDF
Vorticity distribution yields pdf for a single structure
- Evolution in time
Equations for gaussian widths in radius and axial direction and for circulation interms of local strainrate field, longtime survival possible